

$\bar{\theta}, \dot{\bar{\theta}}, \ddot{\bar{\theta}} \longleftrightarrow \bar{\tau} \leftarrow \begin{matrix} \text{Forces + Torques} \\ \text{on each joint} \end{matrix} \leftarrow \begin{matrix} \text{Energy} \\ \text{balance} \end{matrix}$

For control + simulation

Lagrange-Euler equations for robot dynamics

$$\frac{d}{dt} \left[\frac{\partial L}{\partial \dot{q}_i} \right] - \frac{\partial L}{\partial q_i} = \tau_i \quad i=1, 2, \dots, n$$

$L =$ Lagrangian function = Kinetic energy K - Potential energy P

$K =$ Total kinetic energy of robot = $\sum_k dk = \frac{1}{2} \sum_{i=1}^n \sum_{p=1}^3 \sum_{r=1}^3 [\text{Tr}(\bar{U}_{ip} \bar{J}_i \bar{U}_{ir}^T)] \dot{q}_p \dot{q}_r$

$P =$ Total Potential energy of robot = $\sum_i -m_i \bar{g} \cdot {}^0\bar{r}_i = \sum_i -m_i \bar{g} \cdot ({}^0\bar{A}_i \cdot {}^i\bar{r}_i)$

$q_i =$ generalized coordinates (joint or cartesian)

$\dot{q}_i =$ first time derivative

$\tau_i =$ Generalized force (or Torque) applied at joint i to drive link i

$U_{ij} = {}^0\bar{A}_{j-1} \cdot \frac{\partial}{\partial q_j} \left(\overset{Q_j \text{ (see below)}}{{}^j\bar{A}_j \cdot {}^j\bar{A}_i} \right) \quad j \leq i \quad \text{or} \quad = 0, \quad j > i = \frac{\partial {}^0\bar{A}_i}{\partial q_j}$

$\bar{J}_i =$ 4x4 matrix (computed once for the robot)

$m_i =$ mass of link i

$\bar{g}_i =$ gravity row vector = $(g_x, g_y, g_z, 0) = (0, 0, -|g|, 0)$ for a level system

${}^0\bar{r}_i =$ position of link i with respect to frame 0

${}^0\bar{A}_i =$ position of frame i with respect to frame 0

${}^i\bar{r}_i =$ position of center of mass of link i with respect to frame i

or $\bar{\tau}(t) = \bar{D}(\bar{q}(t)) \ddot{\bar{q}}(t) + \bar{h}(\bar{q}(t), \dot{\bar{q}}(t)) + \bar{c}(\bar{q}(t)) = \text{torque vector } (\tau_1, \tau_2, \dots, \tau_n)^T$

$\bar{D}(\bar{q}(t)) =$ Inertial acceleration matrix (nxn)

$$D_{ik} = \sum_{j=\max(i,k)}^n \text{Tr}(\bar{u}_{jk} \bar{J}_j \bar{u}_{ji})$$

$\ddot{\bar{q}}(t) =$ second time derivative of \bar{q}

$\bar{h}(\bar{q}, \dot{\bar{q}}) =$ Coriolis and centrifugal force vector

$$h_i = \sum_{k=1}^n \sum_{m=1}^n h_{ikm} \dot{q}_k \dot{q}_m$$

$$h_{ikm} = \sum_{j=\max(i,k,m)}^n \text{Tr}(\bar{u}_{jk} \bar{J}_j \bar{u}_{ji})$$

$$\bar{u}_{ijk} = \begin{cases} {}^0\bar{A}_{j-1} \cdot \bar{Q}_j \cdot {}^j\bar{A}_{k-1} \cdot \bar{Q}_k \cdot {}^k\bar{A}_i & i < j < k \\ {}^0\bar{A}_{k-1} \cdot \bar{Q}_k \cdot {}^k\bar{A}_{j-1} \cdot \bar{Q}_j \cdot {}^j\bar{A}_i & i < j \text{ or } i < k \\ 0 & \text{otherwise} \end{cases}$$

$$\bar{Q}_i = \begin{cases} \begin{bmatrix} 0 & -1 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix} & \text{revolute joint} \\ \begin{bmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 0 & 0 \end{bmatrix} & \text{prismatic joint} \end{cases}$$

$\bar{c} =$ gravity loading vector $\Rightarrow c_i = \sum_{j=1}^n (-m_j \bar{g} \bar{u}_{ji} \cdot {}^j\bar{r}_j)$

$$j \leq i$$

$$\begin{aligned} \bar{U}_{ij} &= \frac{\partial {}^0 \bar{A}_i}{\partial q_j} = \frac{\partial}{\partial q_j} {}^0 \bar{A}_1 \bar{A}_2 \bar{A}_3 \dots \bar{A}_i \\ &= \frac{\partial}{\partial q_j} ({}^0 \bar{A}_{j-1} {}^{j-1} \bar{A}_j {}^j \bar{A}_i) \\ &= {}^0 \bar{A}_{j-1} \underbrace{\frac{\partial {}^{j-1} \bar{A}_j}{\partial q_j}}_i {}^j \bar{A}_i \end{aligned}$$

$${}^{j-1} \bar{A}_j = \begin{pmatrix} \cos \theta_j & -\cos d_j \sin \theta_j & \sin d_j \sin \theta_j & a_j \cos \theta_j \\ \sin \theta_j & \cos d_j \cos \theta_j & -\sin d_j \cos \theta_j & a_j \sin \theta_j \\ 0 & \sin d_j & \cos d_j & d_j \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

Rotating joint
 $q_j = \theta_j$

$$\frac{\partial {}^{j-1} \bar{A}_j}{\partial q_j} = \begin{pmatrix} -\sin \theta_j & -\cos d_j \cos \theta_j & \sin d_j \cos \theta_j & -a_j \sin \theta_j \\ \cos \theta_j & -\cos d_j \sin \theta_j & \sin d_j \sin \theta_j & a_j \cos \theta_j \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}$$

$$= \begin{pmatrix} 0 & -1 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix} \begin{matrix} \leftarrow \bar{Q}_j \\ {}^{j-1} \bar{A}_j \end{matrix} = \bar{Q}_j {}^{j-1} \bar{A}_j$$

$$\begin{aligned} \frac{\partial {}^0 \bar{A}_i}{\partial q_j} &= {}^0 \bar{A}_{j-1} \bar{Q}_j {}^{j-1} \bar{A}_j {}^j \bar{A}_i \\ &= {}^0 \bar{A}_{j-1} \bar{Q}_j {}^{j-1} \bar{A}_i \end{aligned}$$

$${}^{j-1}\bar{A}_j = \begin{pmatrix} \cos\theta_j & -\cos\alpha_j \sin\theta_j & \sin\alpha_j \sin\theta_j & a_j \cos\theta_j \\ \sin\theta_j & \cos\alpha_j \cos\theta_j & -\sin\alpha_j \cos\theta_j & a_j \sin\theta_j \\ 0 & \sin\alpha_j & \cos\alpha_j & d_j \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

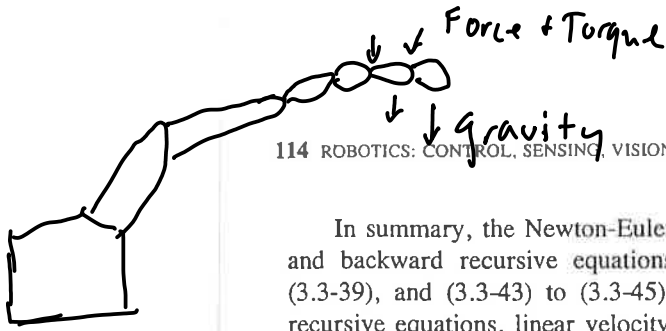
sliding
joint

$$\frac{\partial {}^{j-1}\bar{A}_j}{\partial q_j} = \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 0 & 0 \end{pmatrix} \leftarrow Q_j$$

$$= \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 0 & 0 \end{pmatrix} {}^{j-1}\bar{A}_j$$

Newton-Euler - force balance

many fewer Calculations



In summary, the Newton-Euler equations of motion consist of a set of forward and backward recursive equations. They are Eqs. (3.3-28), (3.3-29), (3.3-35), (3.3-39), and (3.3-43) to (3.3-45) and are listed in Table 3.2. For the forward recursive equations, linear velocity and acceleration, angular velocity and acceleration of each individual link, are propagated from the base reference system to the end-effector. For the backward recursive equations, the torques and forces exerted on each link are computed recursively from the end-effector to the base reference system. Hence, the forward equations propagate kinematics information of each link from the base reference frame to the hand, while the backward equations compute the necessary torques/forces for each joint from the hand to the base reference system.

Table 3.2 Recursive Newton-Euler equations of motion

Forward equations: $i = 1, 2, \dots, n$

$$\bar{\omega}_i = \begin{cases} \bar{\omega}_{i-1} + z_{i-1}\dot{q}_i & \text{if link } i \text{ is rotational} \\ \bar{\omega}_{i-1} & \text{if link } i \text{ is translational} \end{cases}$$

$$\dot{\bar{\omega}}_i = \begin{cases} \dot{\bar{\omega}}_{i-1} + z_{i-1}\ddot{q}_i + \bar{\omega}_{i-1} \times (z_{i-1}\dot{q}_i) & \text{if link } i \text{ is rotational} \\ \dot{\bar{\omega}}_{i-1} & \text{if link } i \text{ is translational} \end{cases}$$

$$\bar{v}_i = \begin{cases} \dot{\bar{\omega}}_i \times p_i^* + \bar{\omega}_i \times (\bar{\omega}_i \times p_i^*) + \bar{v}_{i-1} & \text{if link } i \text{ is rotational} \\ z_{i-1}\dot{q}_i + \dot{\bar{\omega}}_i \times p_i^* + 2\bar{\omega}_i \times (z_{i-1}\dot{q}_i) + \bar{\omega}_i \times (\bar{\omega}_i \times p_i^*) + \bar{v}_{i-1} & \text{if link } i \text{ is translational} \end{cases}$$

$$\bar{a}_i = \dot{\bar{\omega}}_i \times \bar{s}_i + \bar{\omega}_i \times (\bar{\omega}_i \times \bar{s}_i) + \dot{\bar{v}}_i$$

$v_i + \omega_i =$ linear + angular velocities of frame i w.r.t. base. centers of mass or linear acceleration

Backward equations: $i = n, n-1, \dots, 1$

$$F_i = m_i \bar{a}_i$$

$$N_i = I_i \dot{\bar{\omega}}_i + \bar{\omega}_i \times (I_i \bar{\omega}_i)$$

$$f_i = F_i + f_{i+1}$$

$$n_i = n_{i+1} + p_i^* \times f_{i+1} + (p_i^* + \bar{s}_i) \times F_i + N_i$$

3.3-37

$$\bar{v}_i = \begin{cases} \bar{\omega}_i \times \bar{p}_i^* + \bar{v}_{i-1} & \text{rotational} \\ \bar{z}_{i-1}\dot{q}_i + \bar{\omega}_i \times \bar{p}_i^* + \bar{v}_{i-1} & \text{translational} \end{cases}$$

$$\tau_i = \begin{cases} n_i^T z_{i-1} + b_i \dot{q}_i & \text{if link } i \text{ is rotational} \\ f_i^T z_{i-1} + b_i \dot{q}_i & \text{if link } i \text{ is translational} \end{cases}$$

where b_i is the viscous damping coefficient for joint i .

The "usual" initial conditions are $\omega_0 = \dot{\omega}_0 = v_0 = 0$ and $\dot{v}_0 = (g_x, g_y, g_z)^T$ (to include gravity), where $|g| = 9.8062 \text{ m/s}^2$.

① work from base to ee & solve for position, velocity, acceleration
 ② work back from ee to base & solve for forces & torques

calculation

Lagrange-Euler

101,348 * 77,405 +

Newton-Euler

792 * 662 +